Nonlinear Viscoelasticity SINGLE STEP SHEAR STRAIN

$$G(t,\gamma) = \frac{\sigma(t,\gamma)}{\gamma}$$
(5-1)

FINITE RISE TIME



Figure 1: Comparison of Ideal Step Strain (1) with Real Strain Profile (2).

The real strain profile is approximated by a ramp.

$$\gamma(t) = \dot{\gamma}t \tag{5-2}$$

Rise Time
$$\Delta t \equiv \frac{\gamma}{\dot{\gamma}}$$
 (5-3)

Rule of Ten $t > 10\Delta t$ to be meaningful

$$G(t,\gamma) = \sum_{i} G_{i}(\gamma) \exp\left[\frac{-t}{\lambda_{i}(\gamma)}\right]$$
(5-7)

$$G(t,\gamma) = \sum_{i} h_i(\gamma) G_i \exp\left[\frac{-t}{\lambda_i(\gamma)}\right]$$
(5-8)



Figure 2: First Two Moduli (G_i) and Relaxation Times (λ_i) for the Nonlinear Stress Relaxation Modulus of a Polystyrene Solution.

$$G(t,\gamma) = \sum_{i} h_i(\gamma) G_i \exp\left[\frac{-t}{\lambda_i}\right]$$
(5-9)

Nonlinear Viscoelasticity MULTIPLE STEP SHEAR STRAIN



Figure 3: Double Step Strain Experiment.

Boltzmann Superposition works if the steps are small enough to correspond to linear viscoelasticity.

$$\sigma(t) = G(t + t_1)\gamma_1 + G(t)\gamma_2$$
(5-13)

For two nonlinear (large) steps:

$$\sigma(t) = (\gamma_1 + \gamma_2)h(\gamma_1 + \gamma_2)G(t + t_1) + \gamma_2h(\gamma_2)[G(t) - G(t + t_1)]$$
(5-14)

Equation (5-14) works fine with the damping function predicted by the tube model, if the second step was in the same direction as the first step.

Nonlinear Viscoelasticity MULTIPLE STEP SHEAR STRAIN



Figure 4: Double Step Strain Experiment with Reversal.

The double step strain with reversal is a simple experiment that **all** theories of nonlinear viscoelasticity fail to predict.



Figure 5: Spike Strain Test.

Boltzmann superposition predicts no effect of the spike, but experimentally there is an effect when γ_1 is large enough.

Nonlinear Viscoelasticity START-UP OF STEADY SHEAR

Shear Stress Growth Function $\sigma^+(t, \dot{\gamma}) \equiv \sigma(t, \dot{\gamma})$ (5-17)

Shear Stress Growth Coefficient $\eta^+(t,\dot{\gamma}) \equiv \frac{\sigma^+}{\dot{\gamma}}$ (5-18)

First Normal Stress Growth Function $N_1^+(t,\dot{\gamma}) \equiv \sigma_{11}(t,\dot{\gamma}) - \sigma_{22}(t,\dot{\gamma})$

t

(5-19)

 $\Psi_1^+(t,\dot{\gamma}) \equiv \frac{N_1^+}{\dot{\gamma}^2}$ First Normal Stress Growth Coefficient (5-20)

Linear Viscoelastic Limits:

$$\lim_{\dot{\gamma}\to 0} \left[\eta^+(t,\dot{\gamma})\right] = \eta^+(t) \tag{5-23}$$
$$\lim_{\dot{\gamma}\to 0} \left[N_1^+(t,\dot{\gamma})\right] = 0$$

Long Time Limits:

$$\lim_{t \to \infty} \left[\eta^+(t, \dot{\gamma}) \right] = \eta(\dot{\gamma}) \tag{5-24}$$
$$\lim_{t \to \infty} \left[N_1^+(t, \dot{\gamma}) \right] = N_1(\dot{\gamma}) \tag{5-25}$$

Nonlinear Viscoelasticity START-UP OF STEADY SHEAR



Figure 6: Shear Stress Growth and Normal Stress Growth Coefficients for the Start-Up of Steady Shear of a Polystyrene Solution.

Both functions show stress overshoots that indicate short-time relaxation processes are activated in steady shear.

Nonlinear Viscoelasticity CESSATION OF STEADY SHEAR

Shear Stress Decay Function $\sigma^{-}(t,\dot{\gamma}) \equiv \sigma(t,\dot{\gamma})$ (5-33)

Shear Stress Decay Coefficient $\eta^{-}(t,\dot{\gamma}) \equiv \frac{\sigma^{-}}{\dot{\gamma}}$ (5-34)

First Normal Stress Decay Function $N_1^-(t,\dot{\gamma}) \equiv \sigma_{11}(t,\dot{\gamma}) - \sigma_{22}(t,\dot{\gamma})$

 $N_1^-(t,\dot{\gamma}) \equiv \sigma_{11}(t,\dot{\gamma}) - \sigma_{22}(t,\dot{\gamma})$ (5-35)

First Normal Stress Decay Coefficient $\Psi_1^-(t,\dot{\gamma}) \equiv \frac{N_1^-}{\dot{\gamma}^2}$

 $T(t, \dot{\gamma}) \equiv \frac{N_1^-}{\dot{\gamma}^2}$ (5-36)

Linear Viscoelastic Limits:

$$\begin{split} &\lim_{\dot{\gamma}\to 0}\left[\eta^-(t,\dot{\gamma})\right]=\eta^-(t)\\ &\lim_{\dot{\gamma}\to 0}\left[N_1^-(t,\dot{\gamma})\right]=0 \end{split}$$

Short Time Limits:

$$\lim_{t \to 0} \left[\eta^-(t, \dot{\gamma}) \right] = \eta(\dot{\gamma})$$
$$\lim_{t \to 0} \left[N_1^-(t, \dot{\gamma}) \right] = N_1(\dot{\gamma})$$

Nonlinear Viscoelasticity CESSATION OF STEADY SHEAR



Figure 7: Shear Stress Decay and Normal Stress Decay Coefficients for Cessation of Steady Shear Flow of a Polyisobutylene Solution.

Both functions show stresses decaying **faster** at larger shear rates, consistent with long relaxation modes being replaced by shorter-time relaxation processes that are activated in steady shear.

Nonlinear Viscoelasticity NONLINEAR CREEP

Larger $\sigma \Rightarrow$ Larger $\dot{\gamma} \Rightarrow$ Lower η

 $J(t,\sigma) \geq J(t)$

Larger $\sigma \Rightarrow$ More Dissipation Processes (less stored energy)

$$J_s^0(\sigma) \le J_s^0$$



Figure 8: Creep Compliance at a Linear Viscoelastic Stress σ_1 and two Nonlinear Stresses with $\sigma_3 > \sigma_2 > \sigma_1$.

As stress increases, the viscosity drops and the recoverable strain drops, consistent with large stresses inducing additional dissipation mechanisms.

Nonlinear Viscoelasticity NONLINEAR RECOVERY



Figure 9: Creep and Creep Recovery.



Figure 10: Recoverable Compliance after Creep at Three Stress Levels (Increasing Creep Stress from Top to Bottom).

Recoverable compliance is lower at larger creep stresses because the large stress induces additional short-time relaxation processes, meaning that a smaller fraction of the deformation is stored.

Nonlinear Viscoelasticity RECOIL DURING START-UP OF SHEAR

$$\gamma_r \equiv \gamma(t_0) - \gamma(t) = \gamma_r(t - t_0, t_0, \dot{\gamma}) \tag{5-64}$$



Figure 11: Recoil Part-Way Through Start-Up.



Figure 12: Ultimate Recoil During Start-Up Compared with the Shear and Normal Stress Growth Functions for LDPE.